### Gas Accretion onto Magnetized Neutron Stars Ken OHSUGA

**Outflow** 

accretion disk

- Neutron star

Takahashi+ 2018

### Neutron Stars (Ns) & Black Holes (BHs)

#### NS

-a few Msun / 10km-produced by supernovae

Stellar mass BH -about 10 Msun / 30km -produced by supernovae

#### Supermassive BH

- -106-9 Msun
- -located at the galactic center
- -formation mechanism is unknown

# High-energy phenomena

#### NSs/BHs are thought to be engine Cyg X-1 of powerful compact objects

#### SS433





Amy Mioduszewski Michael Rupen Craig Walker Greg Taylor





### Ultra-luminous X-ray sources (ULXs)



# Central object of ULX ?

#### NS

-a few Msun / 10km-produced by supernovae

Big difficulty

Stellar mass BH -about 10 Msun / 30km

-produced by supernovae

Supermassive BH

- -106-9 Msun
- -located at the galactic center

-formation mechanism is unknown

- NO ULXs are offcenter X-ray sources.

### Accretion Power

Potential energy is converted to radiation energy (and jet energy)





-Gas accretion stops if luminosity is larger than Eddington luminosity because of GRAVITY < RADIATION FORCE.

-Thus, the Eddington luminosity is the upper limit of the luminosity.

Neither stellar mass BH (~10Msun) nor NS (~Msun) cannot explain large luminosity of ULXs (>>10<sup>39</sup>erg/s)

## Pulsed ULXs



#### Inconsistency

- -Pulsed emission can be explained by NS.
- -But NS is rejected if the Eddington luminosity is upper limit of the luminosity.

#### <u>Summary</u>

-By performing numerical simulations, we reveal that the NS luminosity can exceed the Eddington limit. -Accretion onto magnetized NS can explain basic features of ULXs.

## **Basic equations**

We solve a full set of general relativistic radiationmagnetohydrodynamic equations

mass cons. 
$$\partial_t \left(\sqrt{-g}\rho u^t\right) + \partial_i \left(\sqrt{-g}\rho u^i\right) = 0$$
  
Gauss's law  $\partial_i \left(\sqrt{-g}B^i\right) = 0$   
Induction eq.  $\partial_t \left(\sqrt{-g}B^i\right) = -\partial_j \left[\sqrt{-g} \left(b^j u^i - b^i u^j\right)\right]$   
energy momentum  
cons. for MHD  $\partial_t \left(\sqrt{-g}T^t_{\nu}\right) + \partial_i \left(\sqrt{-g}T^i_{\nu}\right) = \sqrt{-g}T^{\kappa}_{\lambda}\Gamma^{\lambda}_{\nu\kappa} + \sqrt{-g}G_{\nu}$   
energy momentum  
cons. for radiation  $\partial_t \left(\sqrt{-g}R^t_{\nu}\right) + \partial_i \left(\sqrt{-g}R^i_{\nu}\right) = \sqrt{-g}R^{\kappa}_{\lambda}\Gamma^{\lambda}_{\nu\kappa} - \sqrt{-g}G_{\nu}$   
Radiation

e

# Non-magnetized NS



### Why is super-Eddington feasible?

Mass density



Radiation energy mainly goes in the vertical direction. →Radiation force deos not prevent disk accretion.

Ohsuga et al. 2009 Ohsuga & Mineshige 2011 see also Ohsuga et al. 2005

### Magnetorotational instability



Angular momentum transport induced by magnetorotational instability (MRI) leads to mass accretion

> Hawley & Balbus 1991 See also Velikhov 1959



### Magnetized NS



## Sim. of Accretion column

Radiation energy

NS

x [km]

 $z \, [\rm km]$ 



simulation box

Kawashima et al. 2016



### Model for ULX Pulsar



### SUMMARY

- We performed general relativistic radiation magnetohydrodynamics simulations of gas accretion and outflow around NSs.
- Our simulation revealed that the luminosity can exceed the Eddington luminosity (super-Eddington flow is feasible).
- The central objects of pulsed ULXs would be super-Eddington flows onto the magnetized NSs.