Lattice QCD and nucleon physics: selected excerpts

TSUKUBA GLOBAL SCIENCE WEEK 🚥

TGSW2018 supports the Sustainable Development Goals Mariane Mangin-Brinet



Laboratoire de Physique Subatomique et de Cosmologie

Tsukuba Global Science Week – 21.09.2018



Driving Sustainable elopme,



Introduction

- What is lattice QCD
- > Why do we need it?
- How does it work?
- > How is it organized?
- Reaching the physical pion mass

> Nucleon physics

- nucleon mass
- > axial charge
- > Momentum fraction $\langle x \rangle_{u-d}$ and helicity $\langle x \rangle_{\Delta u-\Delta d}$

Renormalisation

Conclusion

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Prior warning (and apologizes...):

 \geq Presentation limited to what is studied in L.P.S.C. \rightarrow necessarily restricted compared to the richness of the field

Many subjects of interest in LQCD will not be mentioned (nuclear interactions, finite temperature, flavor physics,)

Work performed in the framework of the European Twisted Mass Collaboration (ETMC). Many other competitive international collaborations exist (PACS-CS, MILC, HPQCD, BMW, ALPHA)

What is Lattice QCD?



Courtesy Fermilab Visual Media Services

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What is Lattice QCD?





- Formulated by K. Wilson in 1974 (Phys.Rev. D10 (1974) 2445-2459): gauge theories on a space-time lattice and demonstration that lattice gauge theories at strong coupling lead to confinement
- Regularization of QCD (inherent to all field theories) by space-time discretization
- has become a standard tool

Why do we need it?

- Perturbative QCD (pQCD) very successful in describing many phenomena
- **BUT** limited to an energy range where $\alpha_s \ll 1$

Salient feature of QCD : Asymptotic freedom





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pQCD cannot describe physics at GeV scale

Bound states cannot be described by perturbations

pQCD cannot describe bound states, whatever scale

Understanding bound states in QCD is crucial, because of confinement

Simple example of positronium :

No finite order diagram for $e^+e^- \rightarrow e^+e^-$ has a positronium pole

Positronium pole arises from a divergent infinite sum of diagrams



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Spontaneous breaking of chiral symmetry

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Lattice QCD:

- Only rigorous and systematic non perturbative method to solve QCD ab initio
- Aim: provide solutions of this fundamental theory of matter without uncontrolled uncertainties and with precisions rivaling those of experiments.

How does it work?

- QCD must be regularized in IR and UV
- LQCD is a possible regularization
- Principle:
 - space-time is discretized: $x_{\mu} = a n_{\mu}, \ n_{\mu} \in \mathbb{Z}$
 - quark : $\Psi(x)$ on the lattice sites
 - gluon : $U_{\mu}(x) \in SU(3)$ on the lattice links
 - discretized action: $S = S_G + S_F$



- formulation of QFT in terms of Feynman path integrals with complex time (t \rightarrow it)
- observables obtained by $\langle \mathcal{O} \rangle = \frac{1}{Z} \int DUDq \, D\bar{q} \, \mathcal{O}(U,q,\bar{q}) e^{-S_E}$
- $\langle \mathcal{O} \rangle$ computed by Monte Carlo methods, inspired by statistical physics

Calculations based on stochastic methods: very demanding in terms of computing power.

How it is organized?

- Different QCD action discretizations → different international collaborations
- Pooling of computing resources and manpower inside a given collaboration
- Typical steps of a lattice QCD simulation:
 - 1. Gauge field (or configurations) generation using Monte Carlo methods (HMC) on HPC - most demanding part of the computation
 - 2. Quark propagators calculation by solving a large linear system (HPC)
 - 3. Analysis (or post-treatment), i.e. computations of expectation values of operators, chosen depending on the problem we are interested + statistical study (CPU/GPU farms)



4. Control of systematics (discretization effects,...) by repeating steps 1.2.3. for several lattice parameters (lattice spacing, volume,....)

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Reaching the physical pion mass

Adapted from A. Ukawa, Tsukuba University, « Lattice QCD at the turning point » (2010)



Nucleon physics

Described by an « interpolating field »: $J^N(x) = \varepsilon^{abc} [\mathbf{u}^{Ta}(x) C \gamma_5 \mathbf{d}^b(x)] \mathbf{u}^c(x)$

Observables computed studying the expectation values of operators built from this interpolating field

Simplest example: ab-initio computation of nucleon mass (I)





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Axial charge of the nucleon

• g_A= nucleon axial charge



- Well known from β -decay of the neutron
- Key parameter for understanding the chiral structure of the nucleon
- Considered as a « benchmark » quantity for lattice QCD
- Suffered from a longstanding uncertainty due to its strong dependence in pion mass



 g_A vs pion mass squared for different fermion actions (2010)

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 g_A vs pion mass squared with chiral fits (2010)

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 g_A vs pion mass squared for different fermion actions (2015)

C. Alexandrou et al., Phys.Rev. D92 (2015) 11, 114513, arXiv: 1507.04936

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<u>Momentum fraction $\langle x \rangle_{u-d}$ and helicity $\langle x \rangle_{\Delta u-\Delta d}$ </u>

- An additional step in complexity : 3-point function with derivative operator
- Experimental values extracted from global analysis of collider data
- Has persistently come out too high from lattice calculations
- Concentrated effort to removing each source of systematic errors (pion mass, excited states,...)



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Encouraging recent results → Longstanding puzzle soon resolved ?



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Renormalisation

• Essential ingredient of Lattice QCD calculations

• Discretisation \rightarrow provides a natural regularisation

What is computed on a lattice: « bare » observable: *Obs(a)* (regularized quantities)

• Continuum limit \rightarrow renormalisation

Renormalisation = procedure to take the limit $a \rightarrow 0$: Obs(a) \rightarrow Physical Observables

$$\mathcal{O}_R = Z_\mathcal{O}\mathcal{O}_B$$

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- Renormalization allows, from bare quantities computed on the lattice, to obtain physical observables. It is thus mandatory to make meaningful predictions of observables
- Direct influence of the accuracy of $Z_{\mathcal{O}}$ on observables

Large program of RCs computation using RI-MOM type schemes developed in L.P.S.C. since several years

Dealing with lattice artifacts: an example of challenge to face

Treatment of lattice hypercubic artifacts

- Scalar quantities in the continuum: $Z(p^2)$ rotationally invariant (O(4) group)
- Scalar quantities on an hypercube : $Z(p^2, p^{[4]}, p^{[6]}, p^{[8]}, ...)$ (H(4) group) \rightarrow lattice artifacts
- Illustration see also B. Blossier et al, Phys.Rev. D91 (2015) no.11, 114507)



Quark renormalisation constant Zq vs (ap)² before (blue points) and after (green points) hypercubic artifacts correction Fermion bilinear renormalisation constants (Zq, scalar, pseudo-scalar, axial, vector, tensor) vs (ap)² after hypercubic artifacts correction + running

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- Importance of Nucleon physics for understanding QCD
- > Also crucial for New Physics
- K. Wilson (Lattice 1989, Capri): «One lesson is that LQCD could require a 10⁸ increase of computing power AND spectacular algorithmic advances before ... interaction with experiment takes place ».

At that time, a supercomputer would offer 20 Gflops + « Moore's laws » \rightarrow 10⁸ = 40 years

\rightarrow we are somehow 10 years in advance

- Exciting time for LQCD, with many developments and exp./th. comparisons to come in the next few years
- Developing « strong interactions » between L.P.S.C. and Tsukuba University would be very fruitful

Thank you for your attention

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