

QCD相図探索のための高次ゆらぎ測定と 体積ゆらぎについて

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QCD相転移やQGP生成のモデル化による重イオン衝突の時空発展の
理解に向けた理論・実験共同研究会(Zoom)



筑波大学
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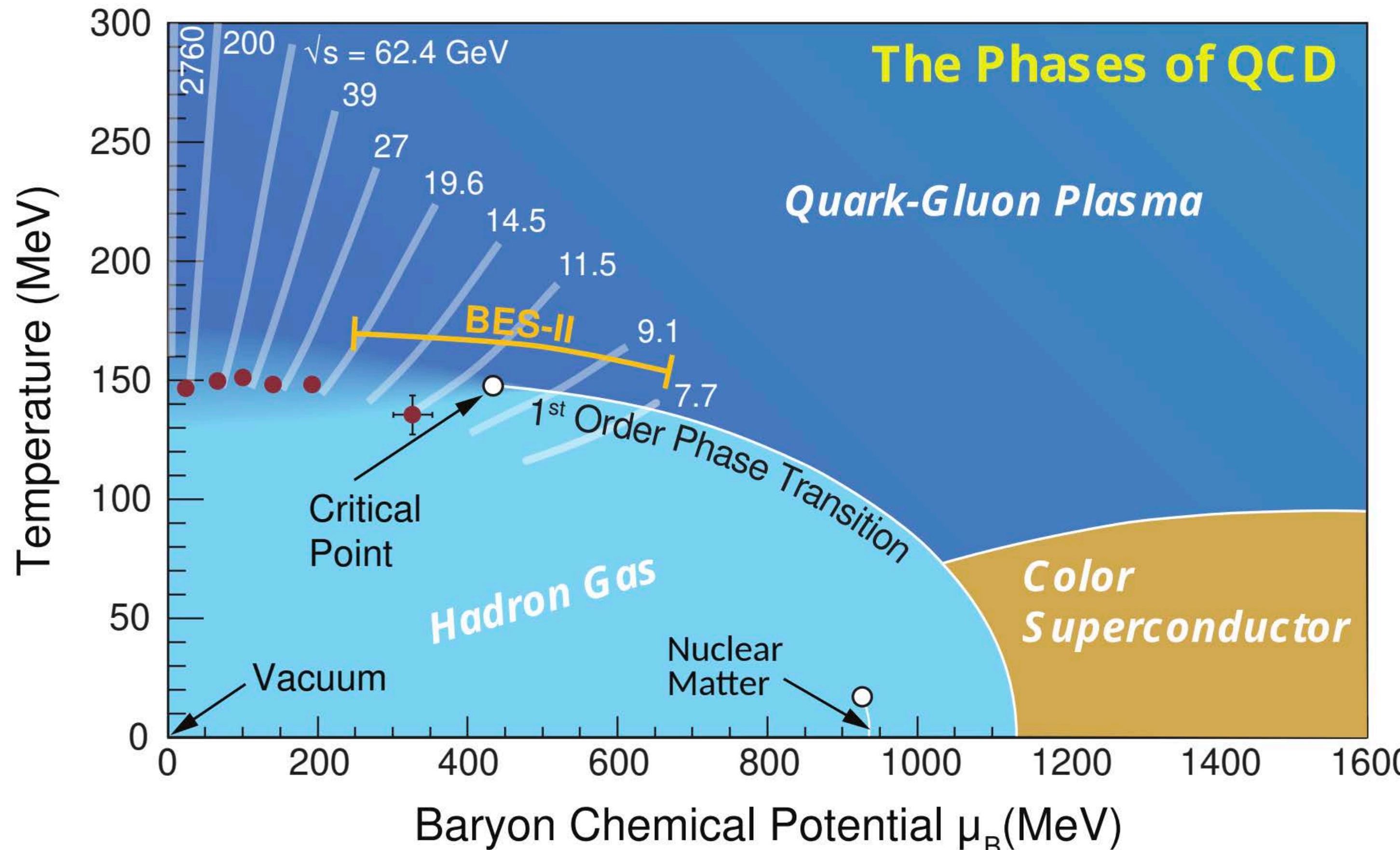
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Outline

- 導入
- 実験結果
- 体積ゆらぎ

QCD phase diagram

✓ QCD phase structure in wide (μ_B, T) region.



- Crossover at $\mu_B = 0$ MeV
Y. Aoki *et al*, *Nature* 443, 675(2006)
- 1st-order phase transition at large μ_B ?
- Critical point?

Beam Energy Scan

✓ Need to investigate the QCD phase structure in wide (μ_B, T) region.

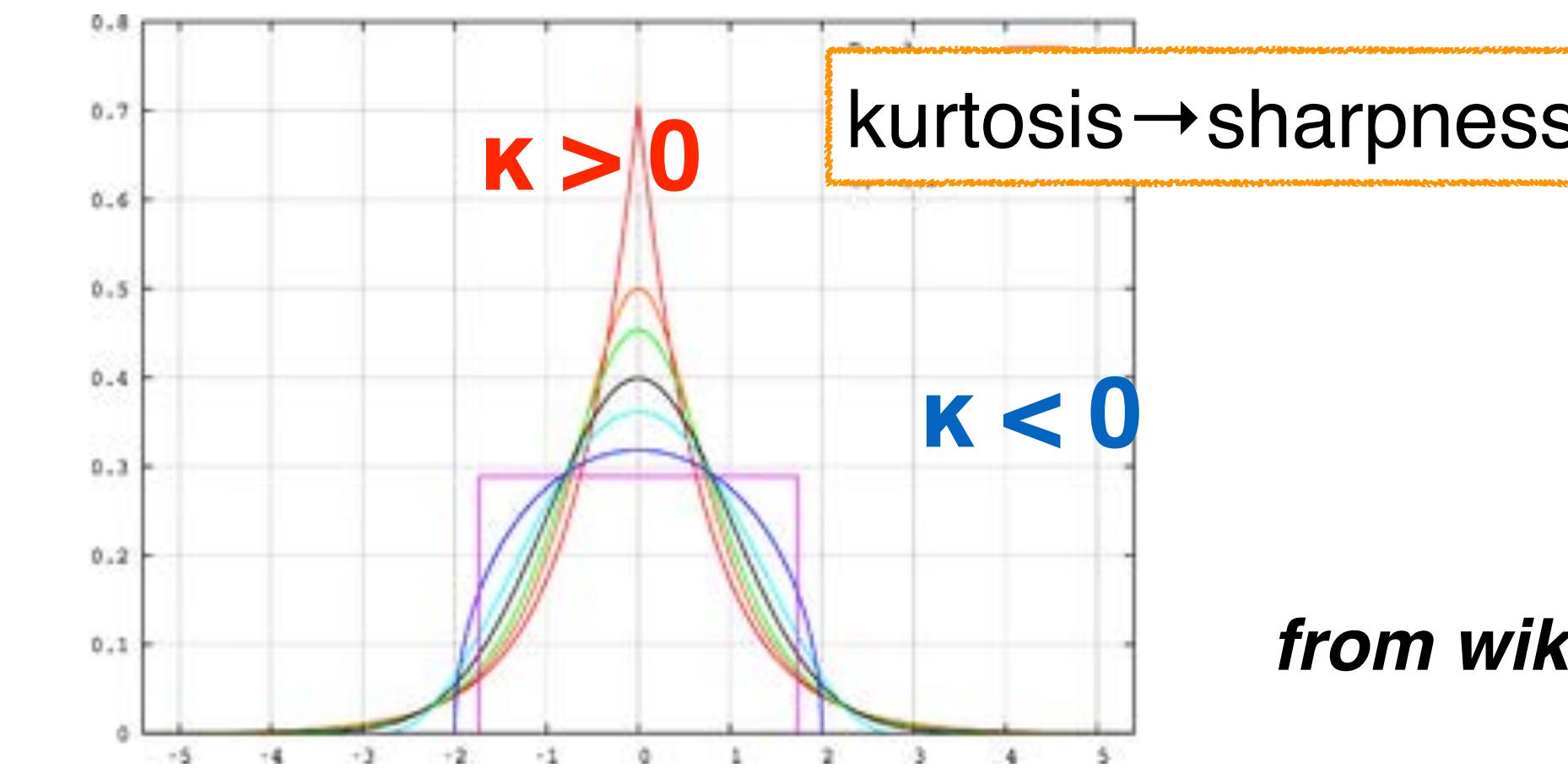
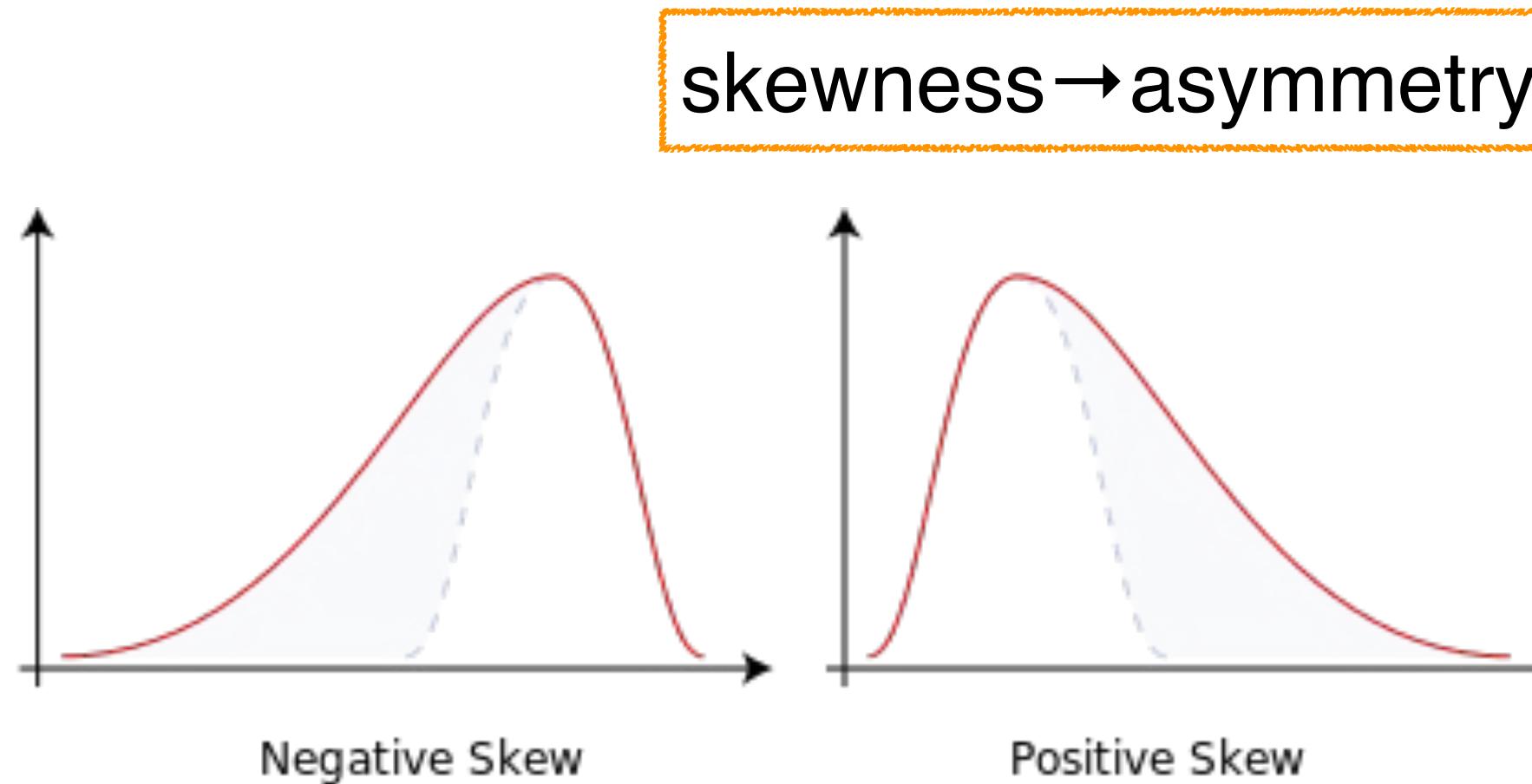
$\sqrt{s_{NN}}$ (GeV)	No. of events (million)	T_{ch} (MeV)	μ_B (MeV)
200	238	164.3	28
62.4	47	160.3	70
54.4	550	160.0	83
39	86	156.4	160
27	2010- 2017	30	155.0
19.6	30	153.9	144
14.5	15	151.6	188
11.5	20	149.4	264
7.7	6.6	144.3	287
	3		398

- Crossover at $\mu_B = 0$ MeV
Y. Aoki et al, Nature 443, 675(2006)
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Higher-order fluctuations

♦ Moments and cumulants are mathematical measures of “shape” of a distribution which probe the fluctuation of observables.

- ✓ Moments: mean (M), standard deviation (σ), skewness (S) and kurtosis (κ).
- ✓ S and κ are sensitive to non-gaussian fluctuations.



from wikipedia

- ✓ Cumulant \Leftrightarrow Central Moment

$$\langle \delta N \rangle = N - \langle N \rangle$$

$$C_1 = M = \langle N \rangle$$

$$C_2 = \sigma^2 = \langle (\delta N)^2 \rangle$$

$$C_3 = S\sigma^3 = \langle (\delta N)^3 \rangle$$

$$C_4 = \kappa\sigma^4 = \langle (\delta N)^4 \rangle - 3\langle (\delta N)^2 \rangle^2$$

- ✓ Cumulant : additivity

$$C_n(X + Y) = C_n(X) + C_n(Y)$$

→ proportional to volume

Fluctuations of conserved quantities

PRL 105, 022302 (2010) :
STAR Collaboration

♦ Net baryon, net charge and net strangeness

“Net” : positive - negative

$$\Delta N_q = N_q - N_{\bar{q}}, \quad q = B, Q, S$$

No. of positively charged particles in one collision

No. of negatively charged particles in one collision

Fill in histograms over many collisions

(1) Sensitive to correlation length

$$C_2 = \langle (\delta N)^2 \rangle_c \approx \xi^2 \quad C_5 = \langle (\delta N)^5 \rangle_c \approx \xi^{9.5}$$

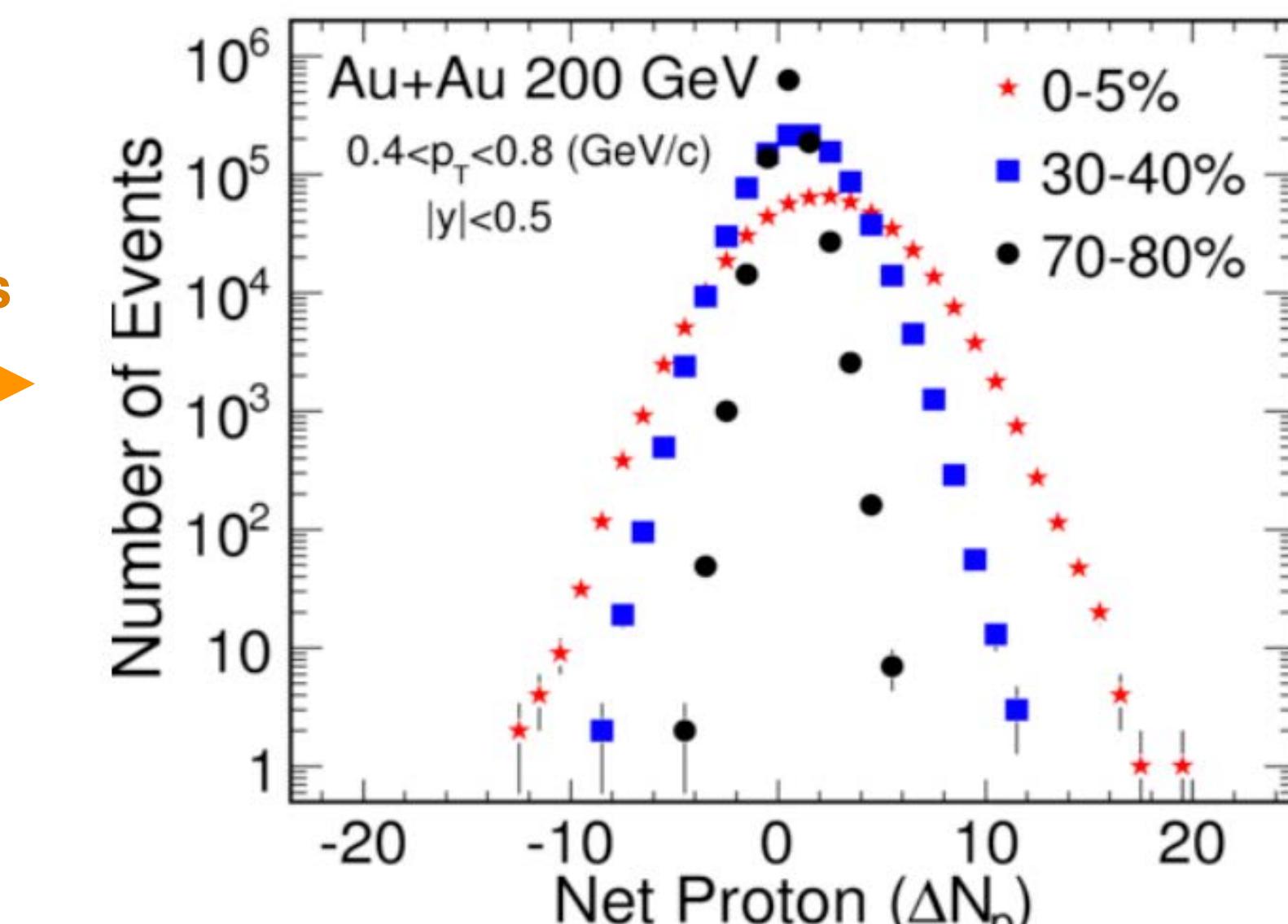
$$C_3 = \langle (\delta N)^3 \rangle_c \approx \xi^{4.5} \quad C_6 = \langle (\delta N)^6 \rangle_c \approx \xi^{12}$$

$$C_4 = \langle (\delta N)^4 \rangle_c \approx \xi^7$$

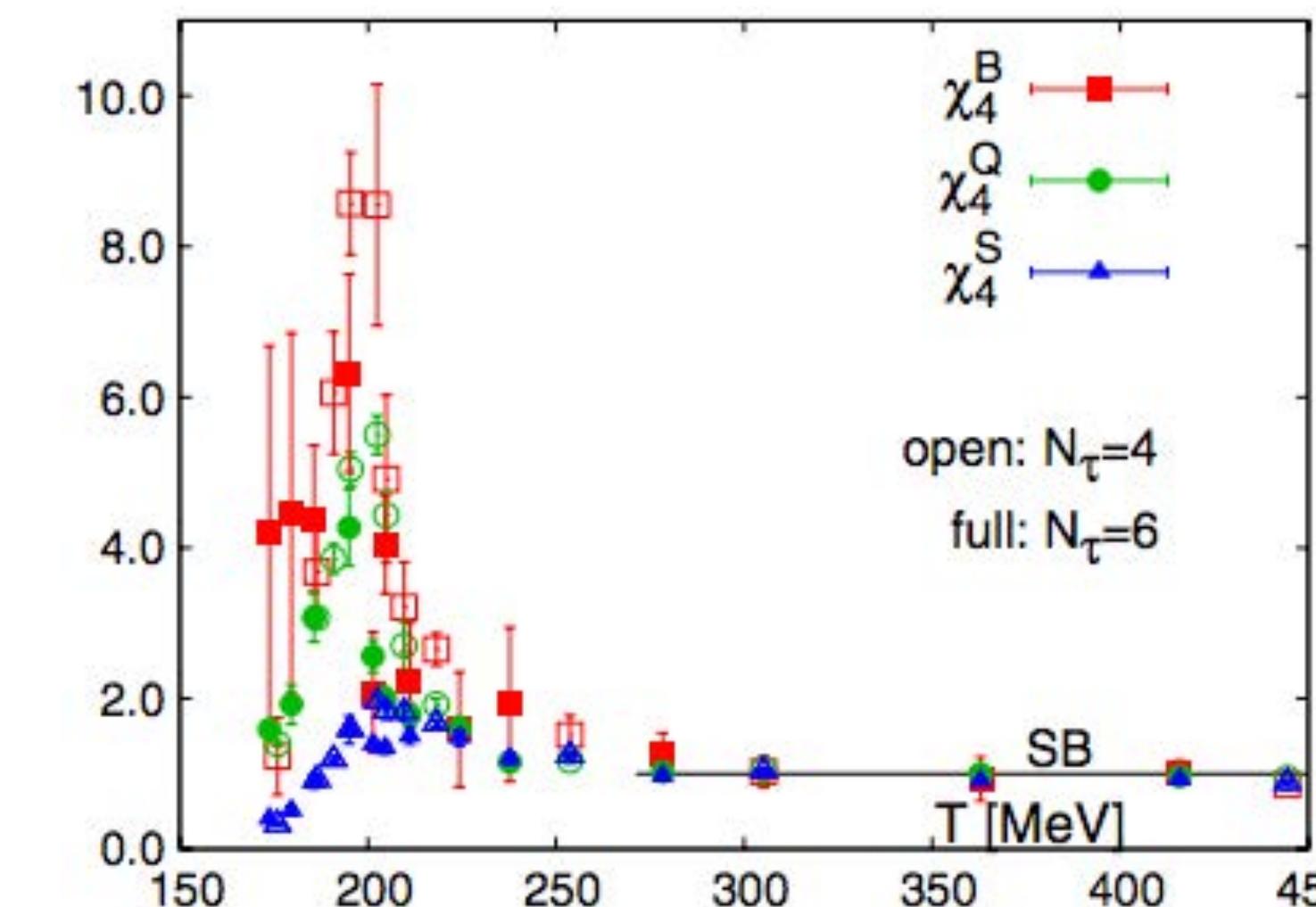
M. A. Stephanov, Phys. Rev. Lett. 102, 032301 (2009)

M. A. Stephanov, Phys. Rev. Lett. 107, 052301 (2011)

MAkawa, S. Ejiri and M. Kitazawa, Phys. Rev. Lett. 103, 262301 (2009)



→neutrons cannot be measured



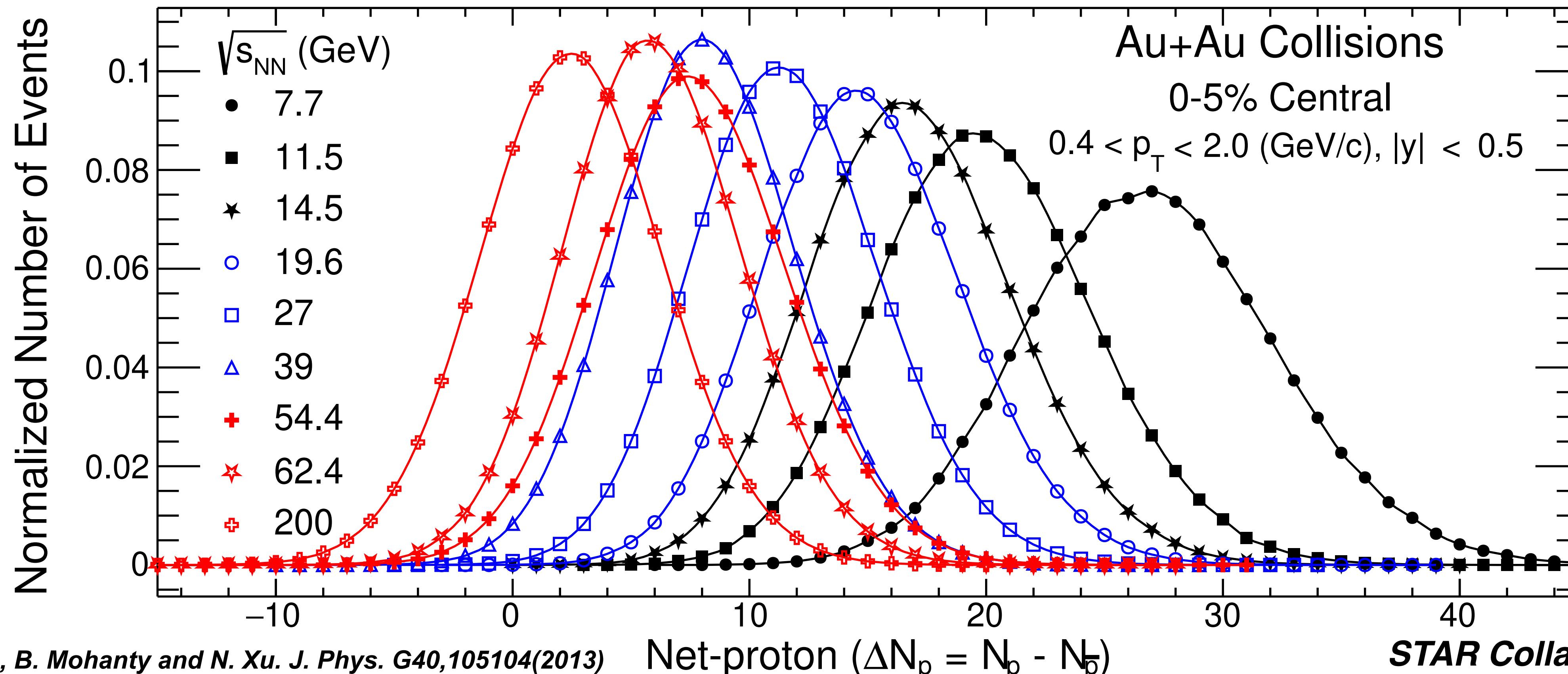
M. Cheng et al, PRD 79, 074505 (2009)

$$S\sigma = \frac{C_3}{C_2} = \frac{\chi_3}{\chi_2} \quad \kappa\sigma^2 = \frac{C_4}{C_2} = \frac{\chi_4}{\chi_2}$$

$$\chi_n^q = \frac{1}{VT^3} \times C_n^q = \frac{\partial^n p/T^4}{\partial \mu_q^n}, \quad q = B, Q, S$$

Raw net-proton distribution

- ✓ Avoid auto-correlation effects : New centrality definition
- ✓ Suppress initial volume fluctuation : Centrality bin width correction
- ✓ Detector efficiency correction : Binomial model



X.Luo, J. Xu, B. Mohanty and N. Xu. J. Phys. G40,105104(2013)

M. Kitazawa : PRC.86.024904(2012)

A. Bzdak and V. Koch : PRC.86.044904(2012), X. Luo : PRC.91.034907(2016)

T. Nonaka, M. Kitazawa, S. Esumi : PRC.95.064912(2017), NIMA906 10-17 (2018),

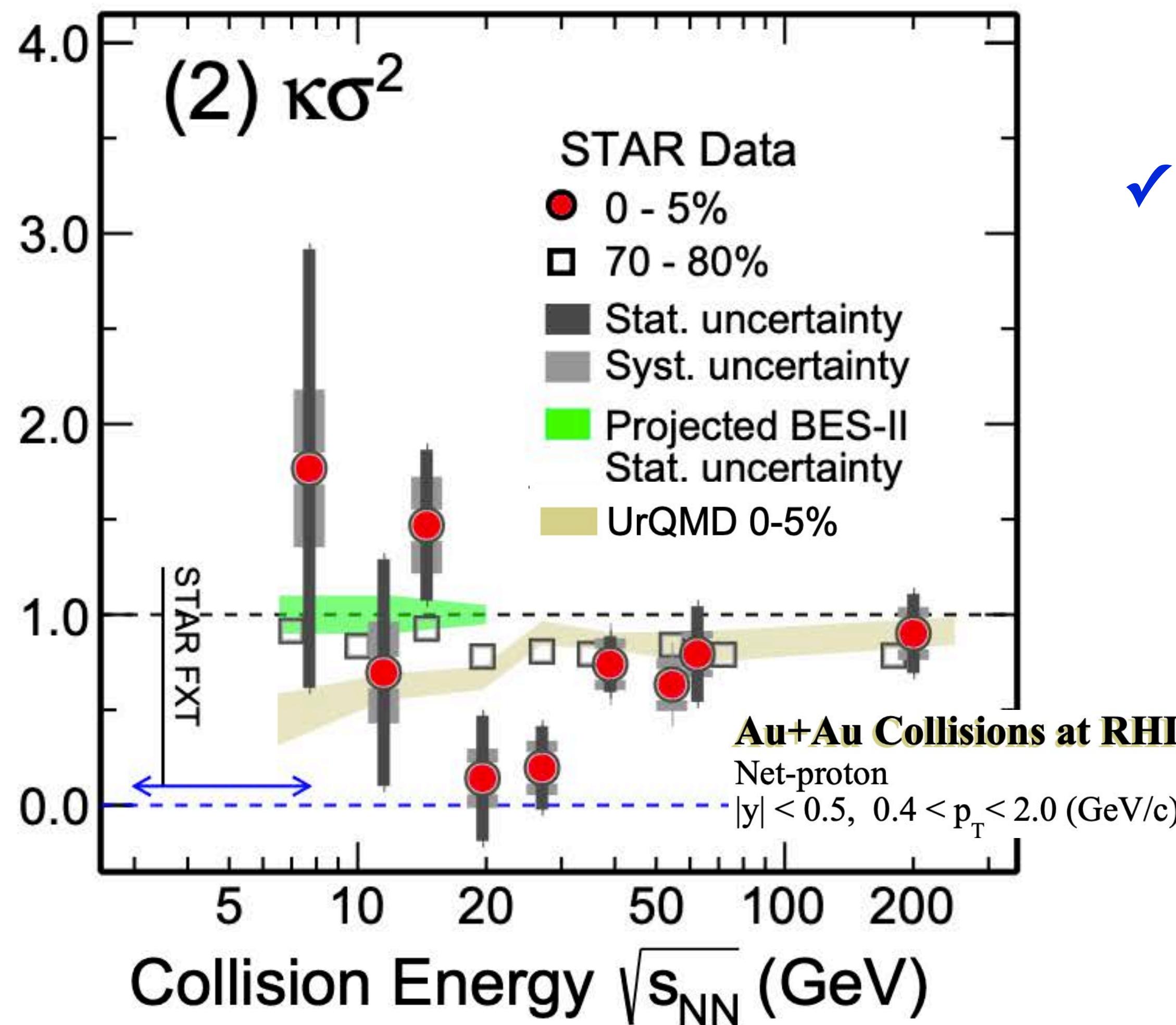
NIMA984(2020)164632

X. Luo, T. Nonaka : PRC.99.044917(2019)

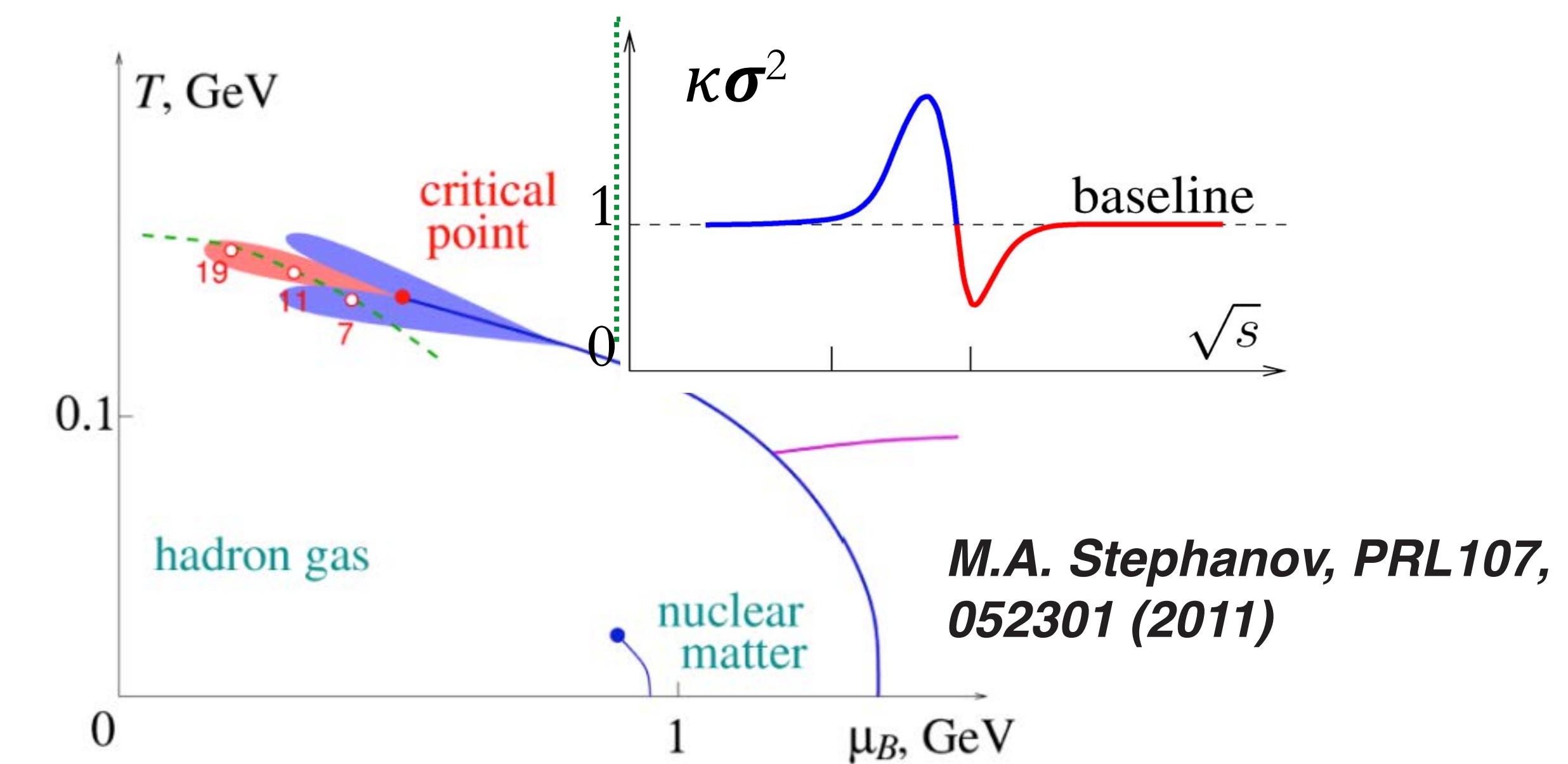
STAR Collaboration,
PRL.126.092301(2021)
PRC.104.024902(2021)

C_4/C_2 for critical point search

STAR Collaboration, PRL.126.092301(2021)

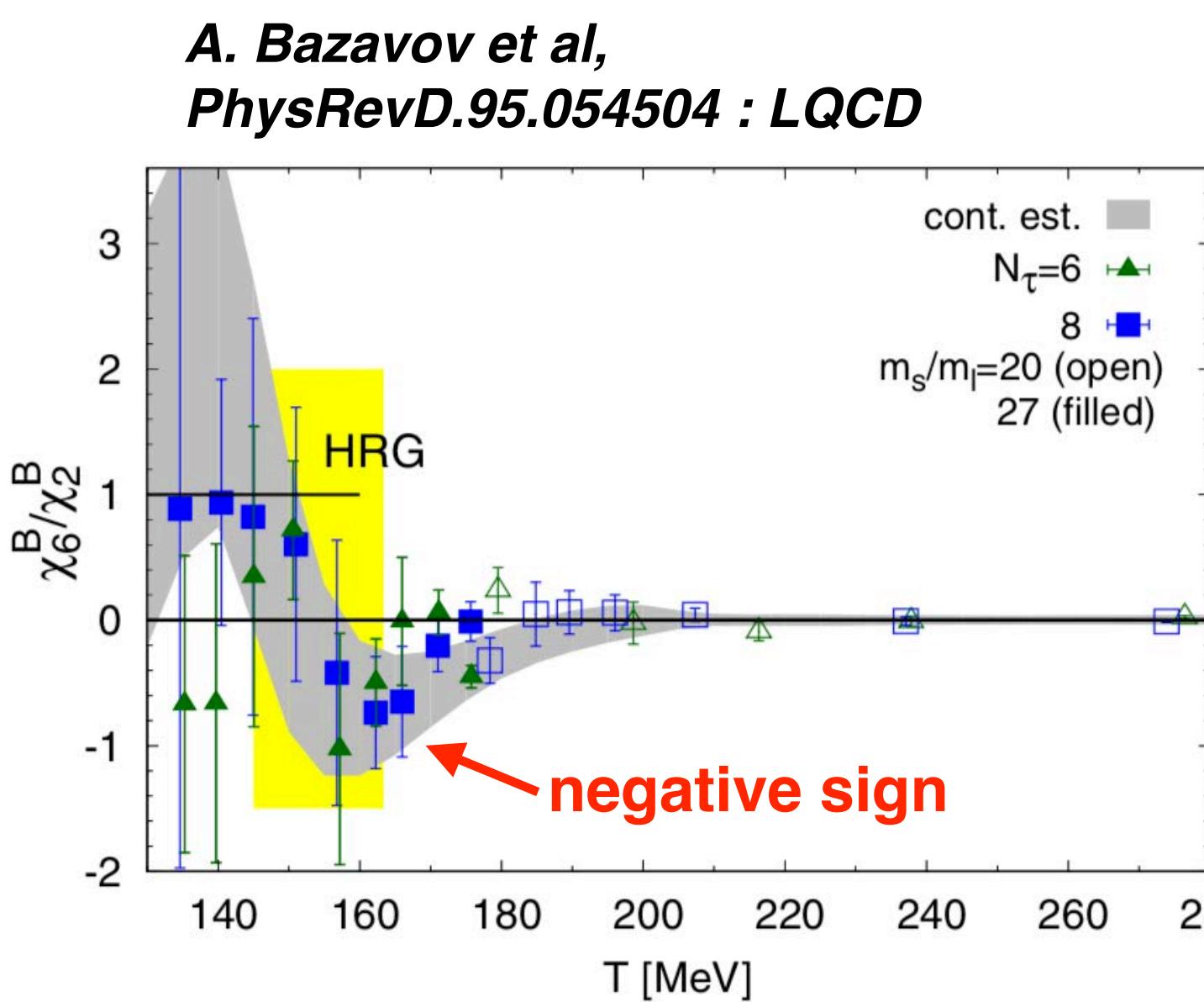


- ✓ Net-proton $\kappa\sigma^2$ (C_4/C_2) shows a non-monotonic behaviour. The trend is consistent with the expectation from theoretical calculations having a critical point.
- ✓ Enhancement at low beam energies cannot be explained by baryon number conservation.

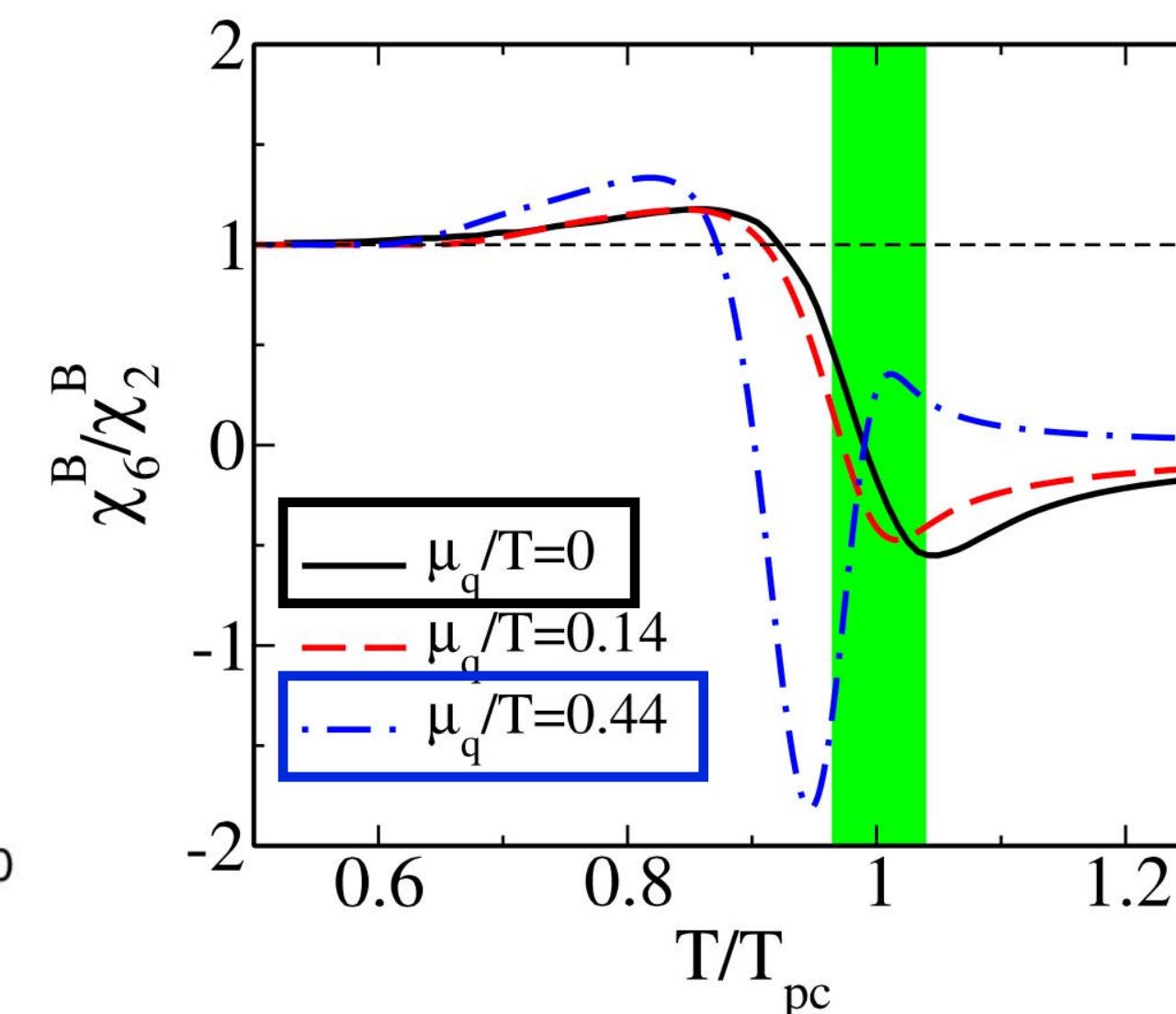


C_6/C_2 for crossover search

- ✓ There isn't yet any direct experimental evidence for the smooth crossover at $\mu_B \sim 0$.
- ✓ $C_6/C_2 < 0$ is predicted as a signature of crossover transition.
- ✓ High-statistics data sets at $\sqrt{s_{NN}} = 27, 54.4, \text{ and } 200 \text{ GeV}$ are analyzed to look for the **experimental signature of crossover transition**.



Friman et al, Eur. Phys. J. C (2011)
71:1694 : PQM model



C.Schmidt, Prog.Theor.Phys.Supp.186,563–566(2010)

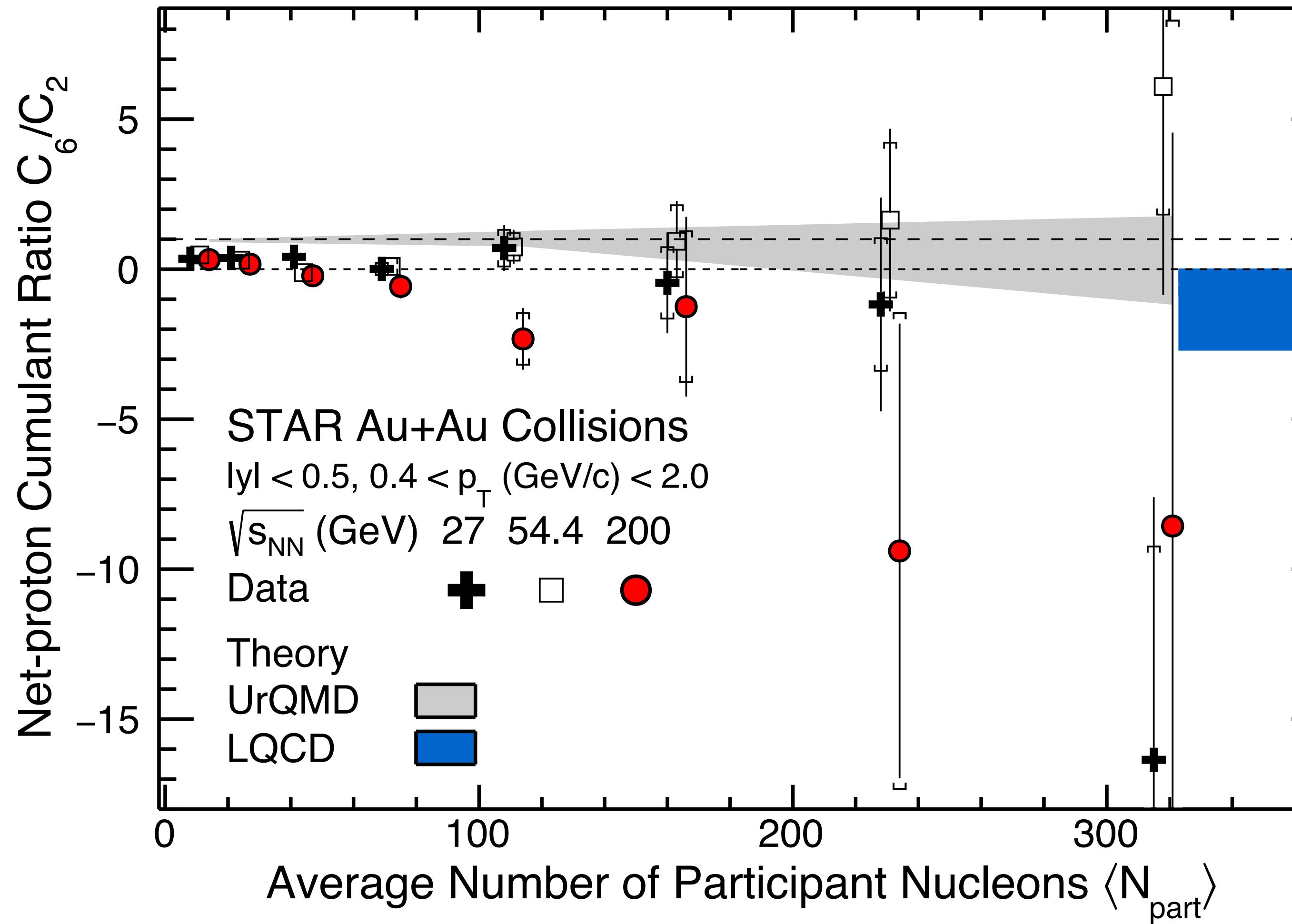
Cheng et al, Phys. Rev. D 79, 074505 (2009)

Friman et al, Eur. Phys. J. C (2011) 71:1694

Freeze-out conditions	χ_4^B/χ_2^B	χ_6^B/χ_2^B	χ_4^Q/χ_2^Q	χ_6^Q/χ_2^Q
HRG	1	1	~2	~10
QCD: $T^{\text{freeze}}/T_{pc} \lesssim 0.9$	$\gtrsim 1$	$\gtrsim 1$	$\gtrsim 1$	~2
QCD: $T^{\text{freeze}}/T_{pc} \simeq 1$	~ 0.5	<0	~1	<0

Predicted scenario for this measurement

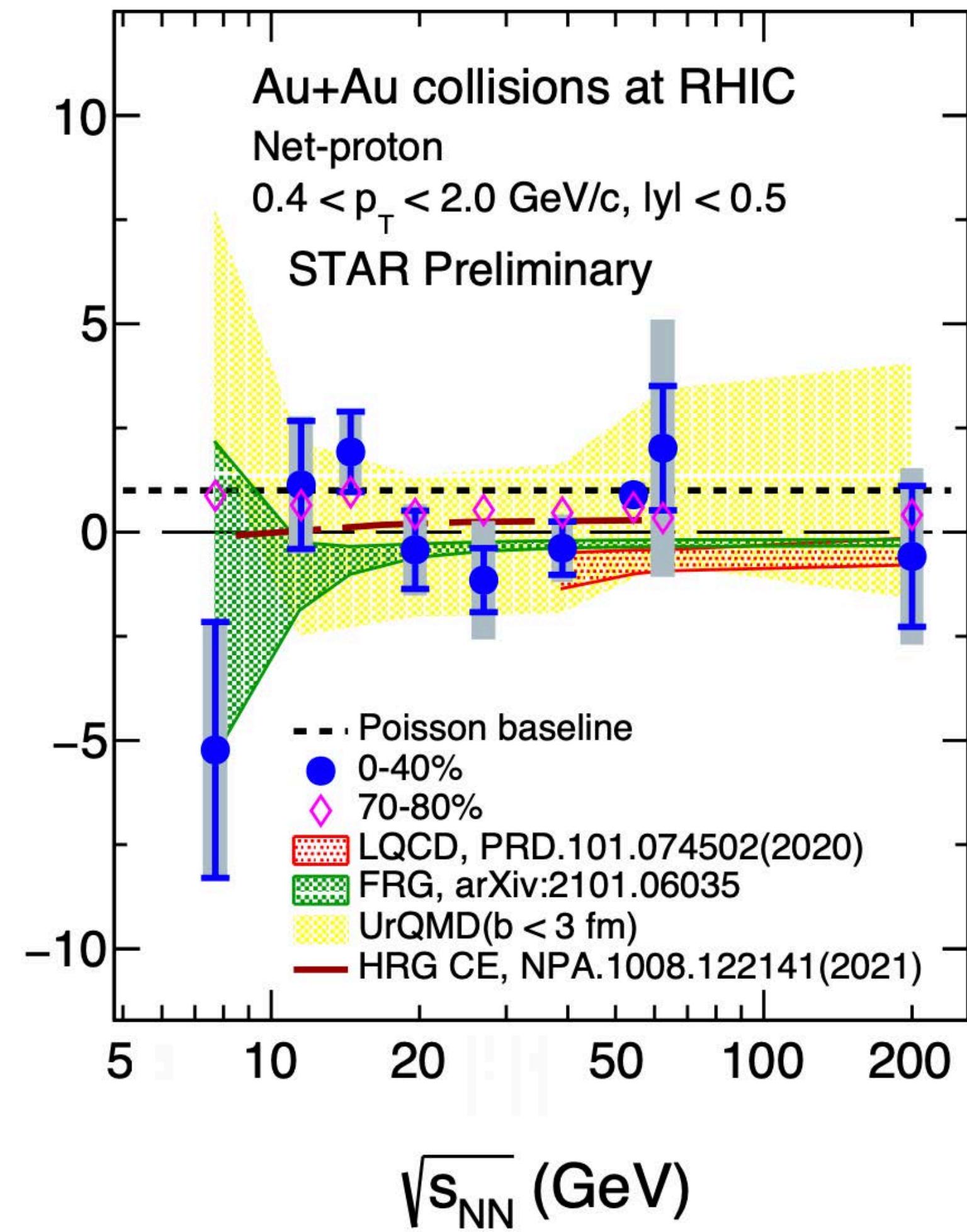
Centrality dependence of C_6/C_2



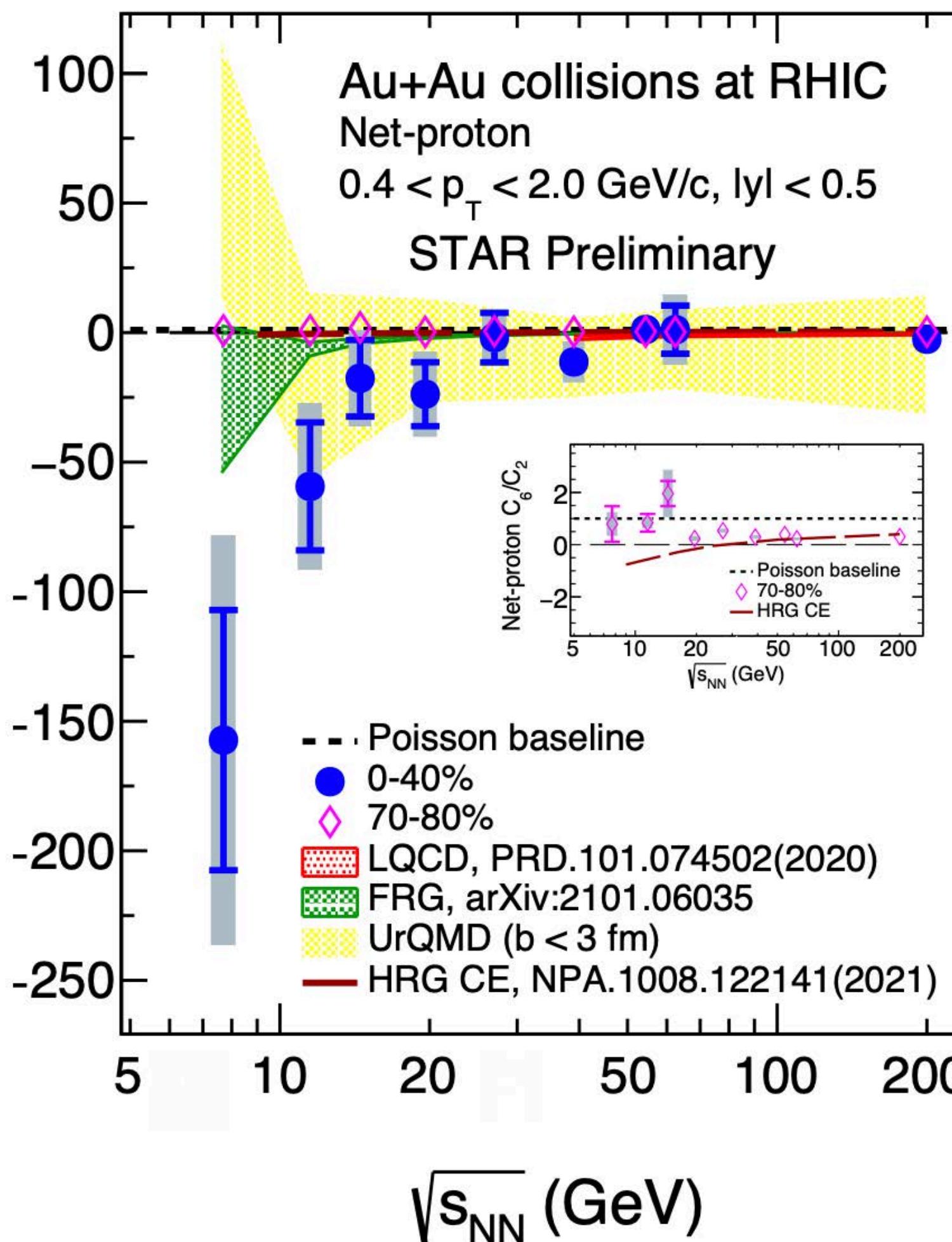
- ✓ C_6/C_2 values are progressively negative from peripheral to central collisions at 200 GeV, which is consistent with LQCD calculations.
- ✓ Could suggest a smooth crossover transition at top RHIC energy.

Energy dependence of C_5/C_1 and C_6/C_2

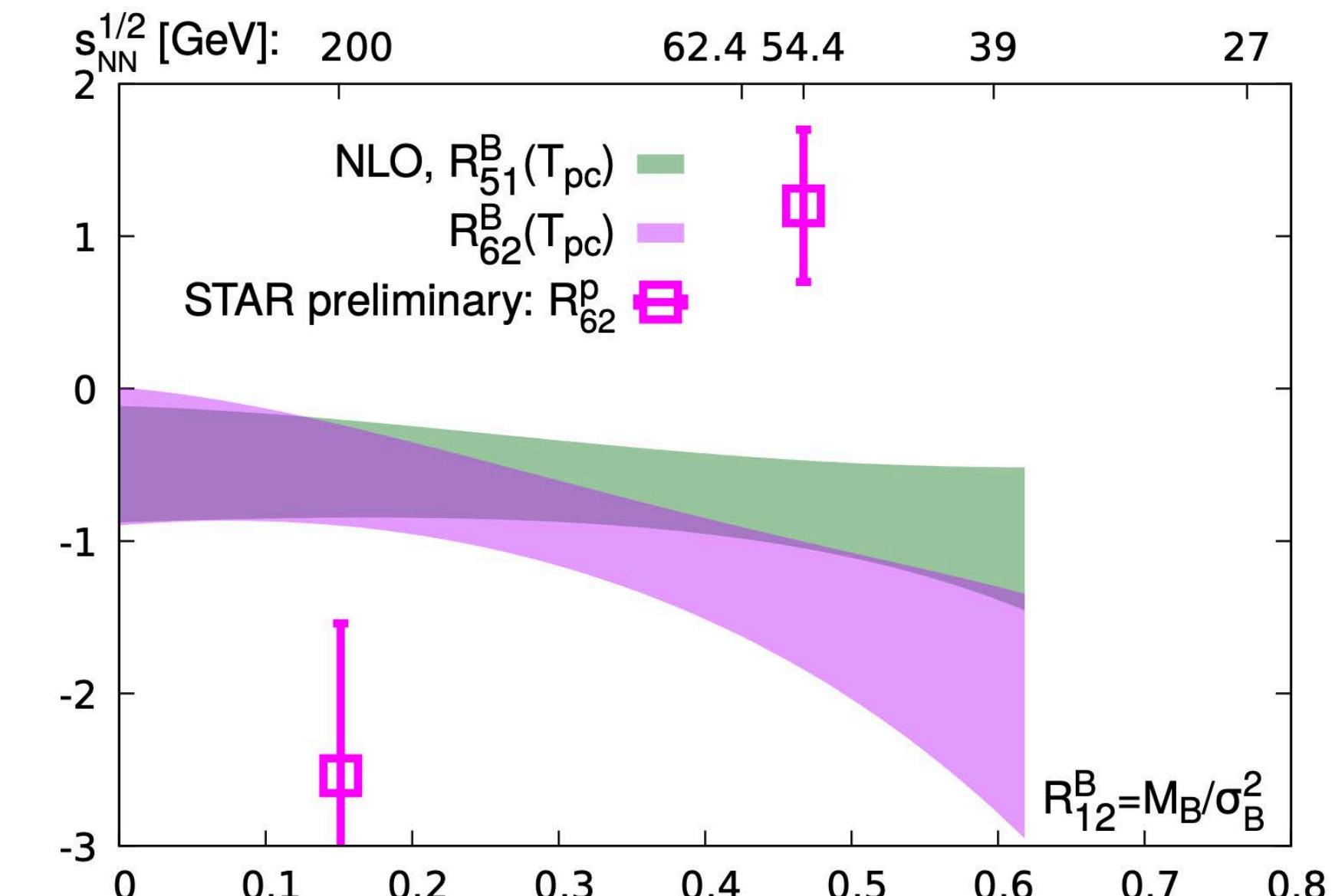
C_5/C_1



C_6/C_2



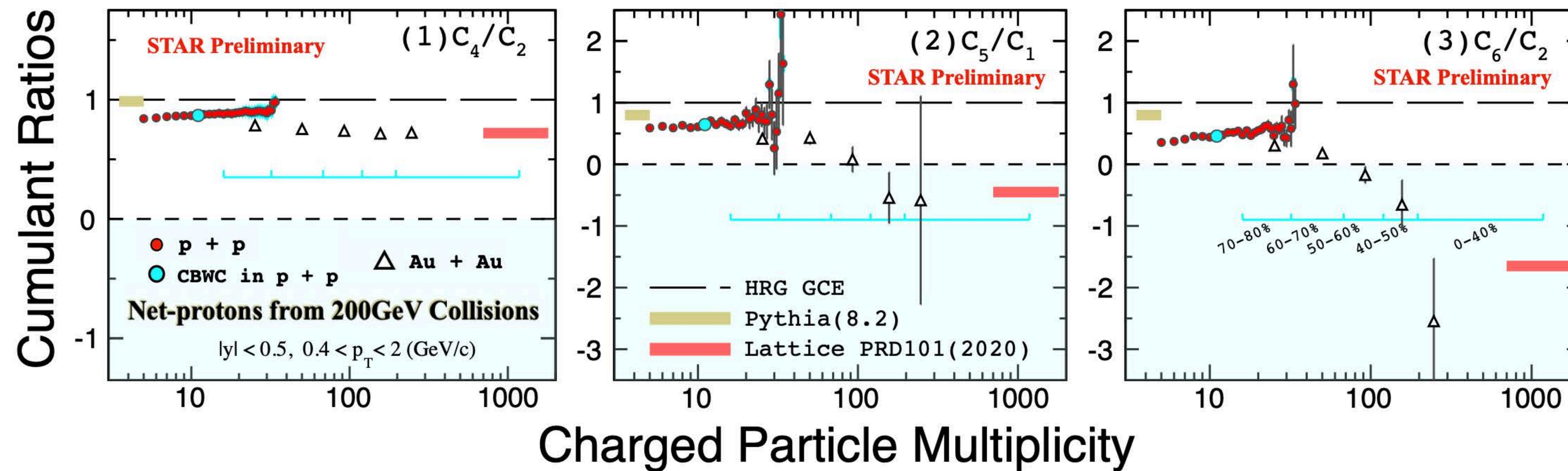
- ✓ Weak collision energy dependence observed for 0-40% centrality.
- ✓ Deviations from zero at a level of $< 2\sigma$ observed for 0-40% centrality.



Bazavov et al., Phys.Rev.D101,074502 (2020)

Multiplicity dependence

- ✓ C_5/C_1 and C_6/C_2 are positive for p+p collisions, while negative for central Au+Au collisions.
- ✓ Lattice calculations imply chiral phase transition in the thermalized QCD matter, which is not the case in 200 GeV p+p collisions.



- Only statistical errors are shown for Au+Au results
- Efficiency is not corrected for x-axis

STAR Collaboration,

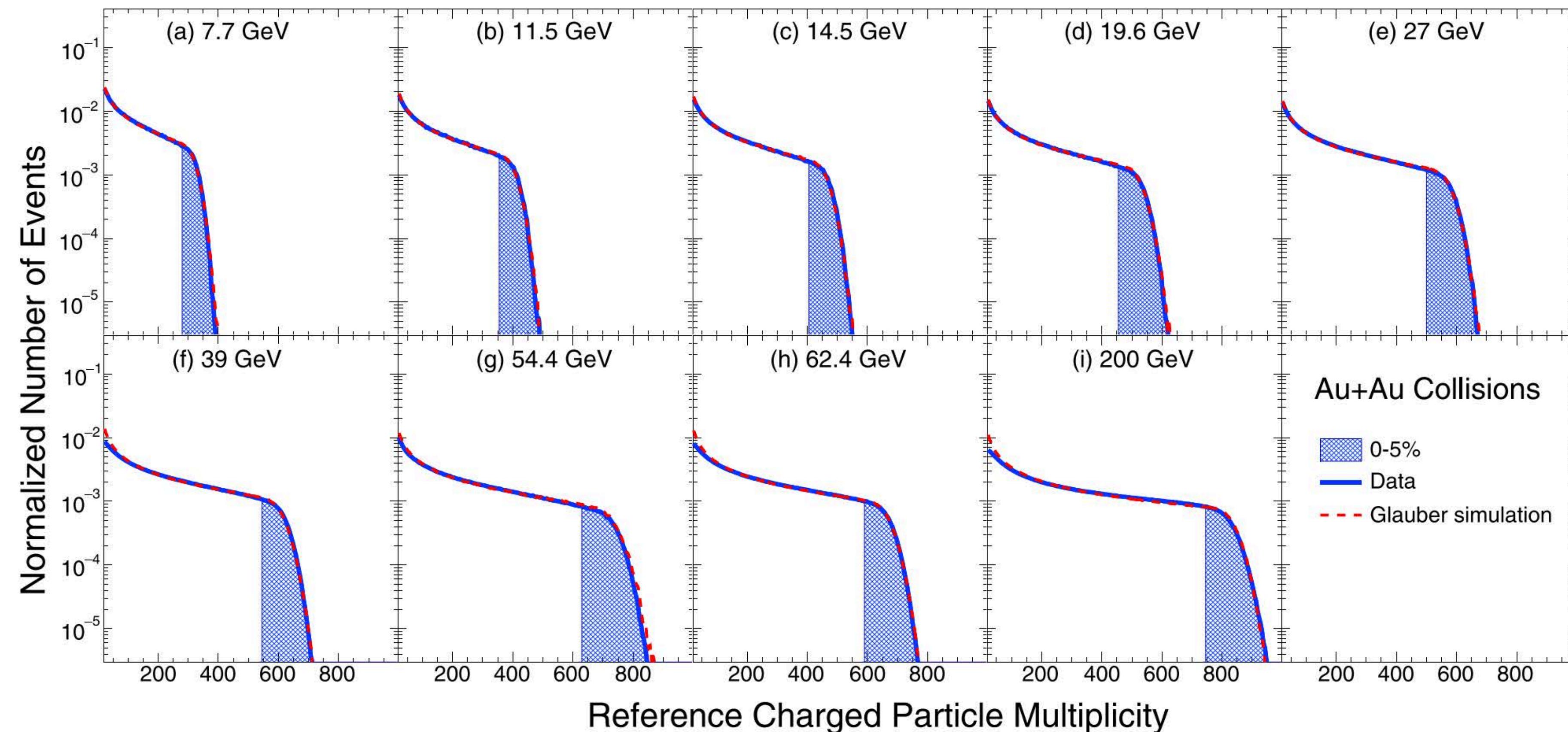
PRC.104.024902(2021)

LQCD : Phys. Rev. D 101, 074502 (2020)

STAR Collaboration,
Nuclear Physics A, 1005,
121882 (2021)

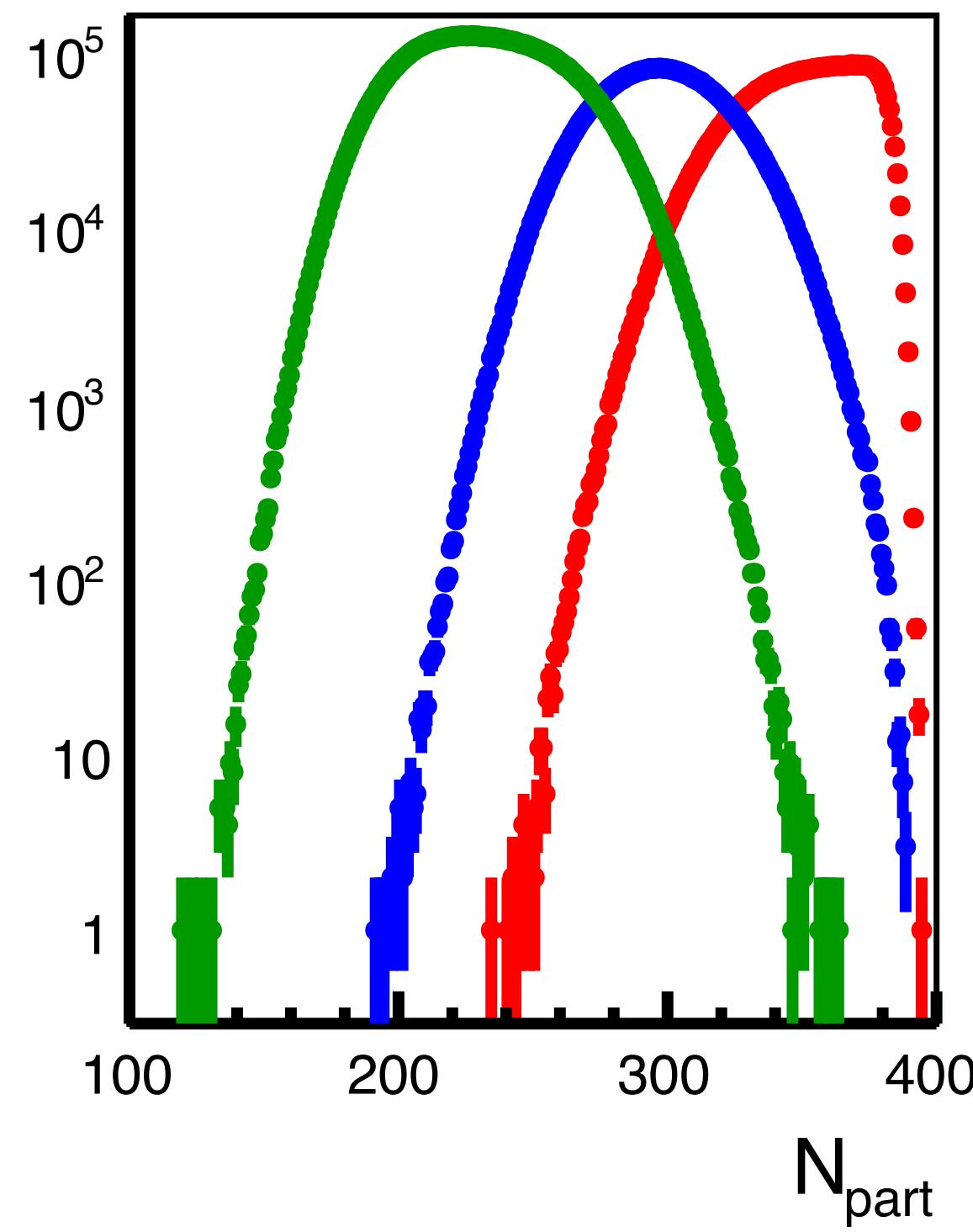
Centrality

✓ インパクトパラメータを測定できないので、粒子数分布（をモデルでフィットした分布）を等分割してCentralityを定義。

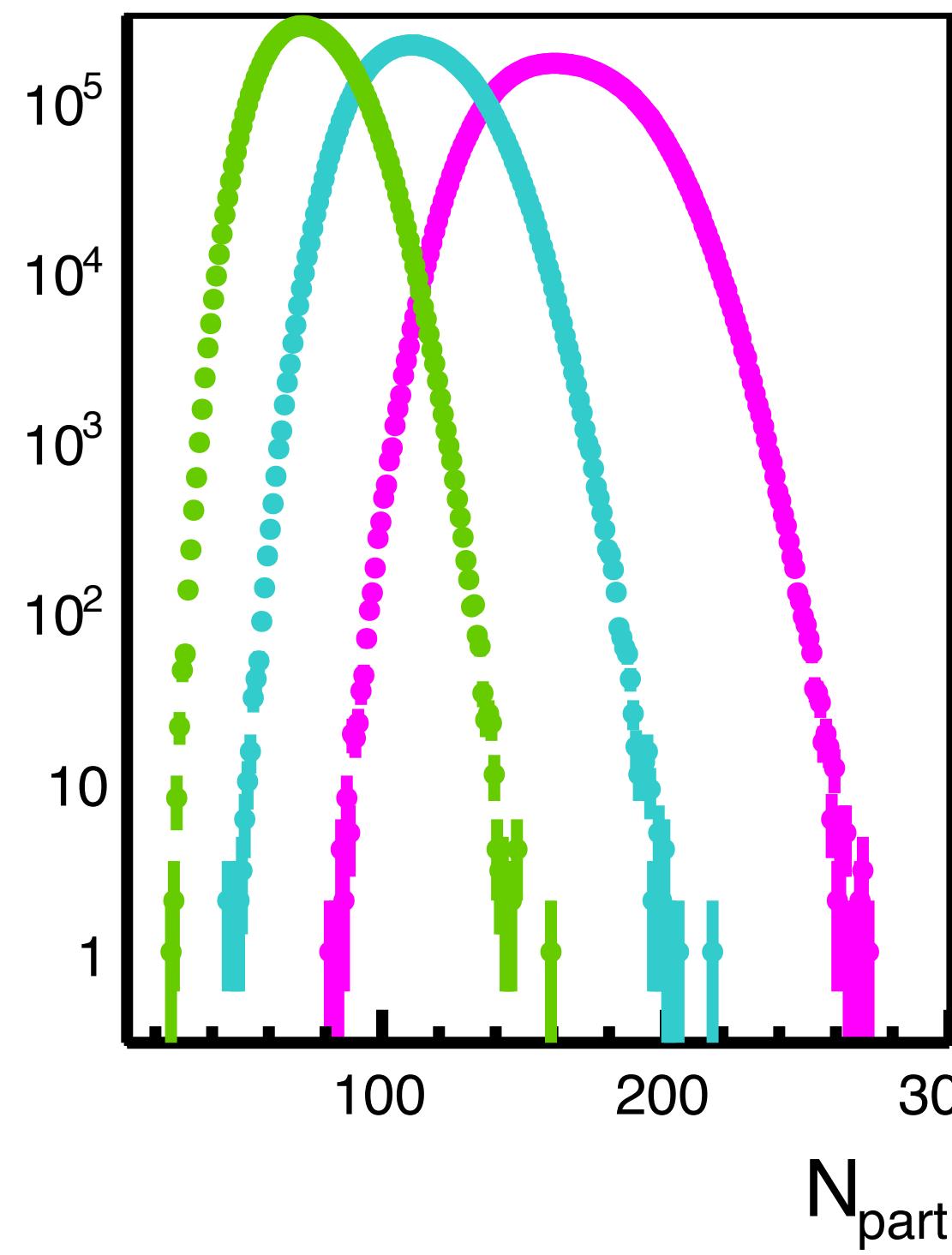


Npart ゆらぎ

- 0-5%
- 5-10%
- 10-20%

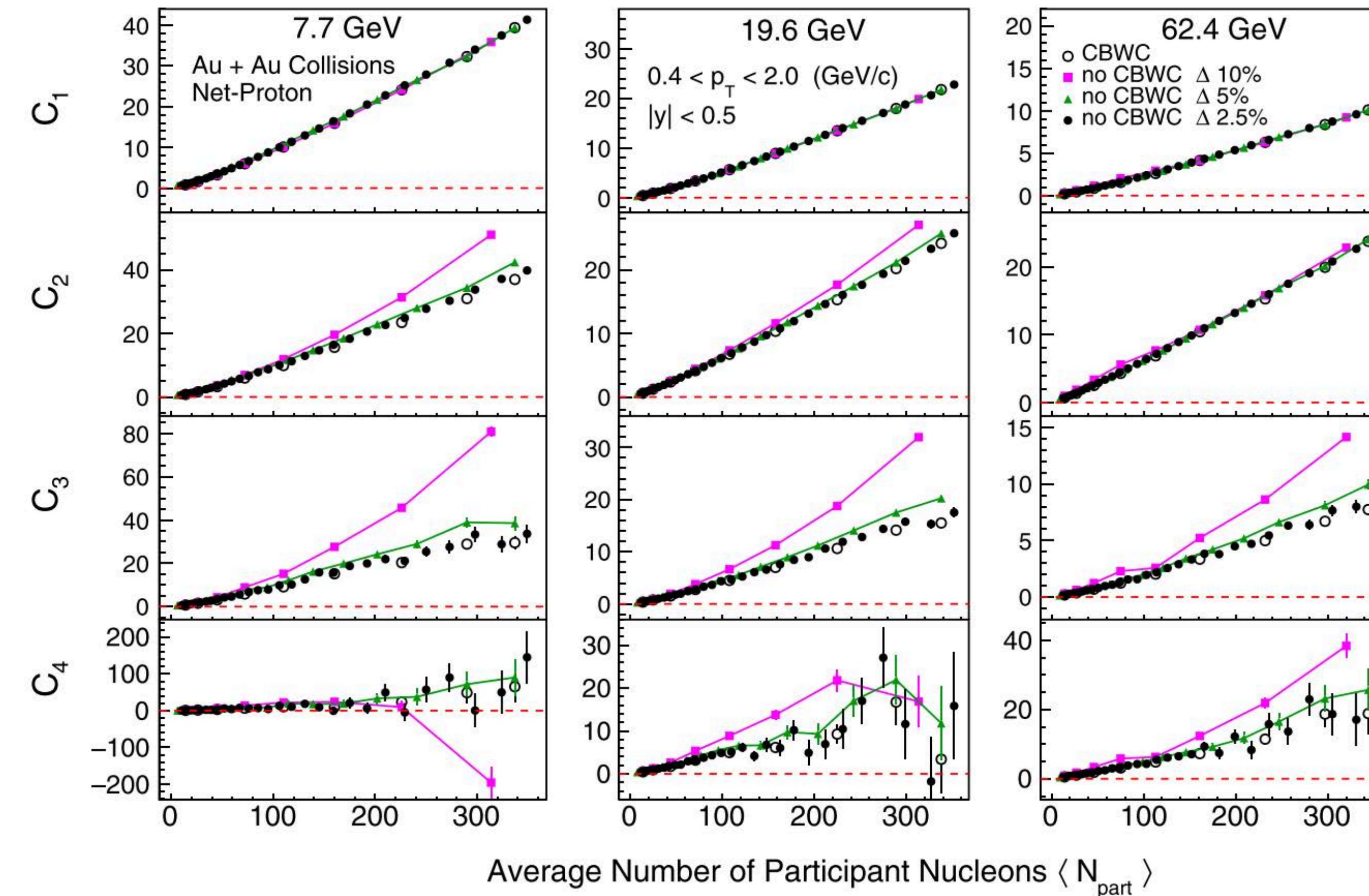


- 20-30%
- 30-40%
- 40-50%



✓ 生成粒子数分布を分割して
Centralityを定義しているため、
 N_{part} (or b)分布は大きく揺らぐ
→体積ゆらぎ

補正手法1 : Centrality bin width correction



- CBWC
- no CBWC $\Delta 10\%$
- ▲ no CBWC $\Delta 5\%$
- no CBWC $\Delta 2.5\%$

- ✓ Centrality幅を狭くすると、ある値に収束→体積ゆらぎの抑制
- ✓ 平均値の解析には影響なし。
- ✓ Reference multiplicity 1 ビンごとにキュムラントを計算し、Centrality内で平均を取る。

補正手法2 : Volume fluctuation correction

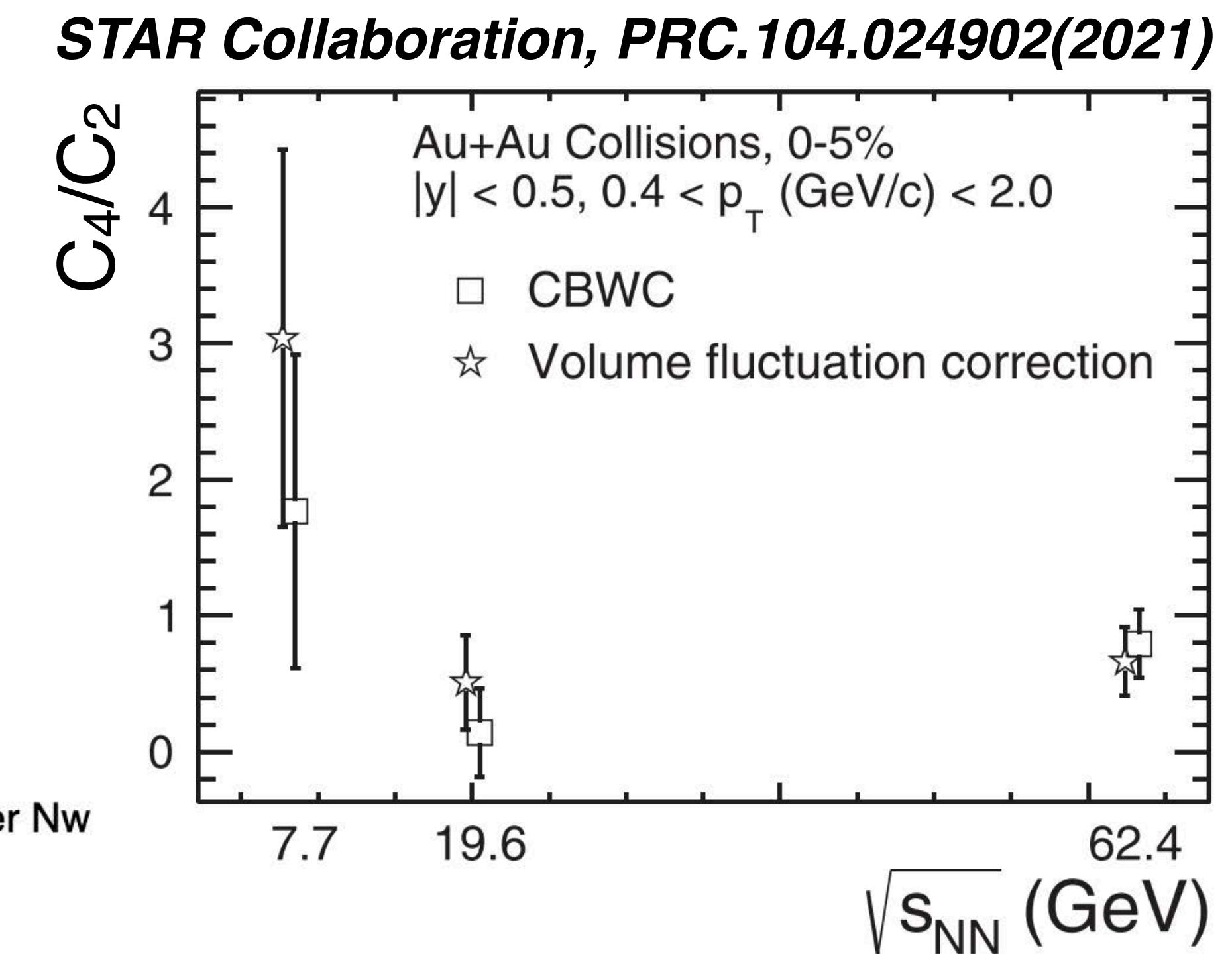
- ✓ Npartごとの独立な粒子生成を仮定すると、測定キュムラントは（真のゆらぎ）と（Npart高次ゆらぎ）の組み合わせで表される。
- ✓ Data-drivenな手法と統計誤差の範囲内で一致。

Measured cumulant	True cumulant
$\kappa_1(\Delta N)$	$= \langle N_W \rangle \kappa_1(\Delta n)$
$\kappa_2(\Delta N)$	$= \langle N_W \rangle \kappa_2(\Delta n) + \langle \Delta n \rangle^2 \kappa_2(N_W),$
$\kappa_3(\Delta N)$	$= \langle N_W \rangle \kappa_3(\Delta n) + 3 \langle \Delta n \rangle \kappa_2(\Delta n) \kappa_2(N_W) + \langle \Delta n \rangle^3 \kappa_3(N_W),$
$\kappa_4(\Delta N)$	$= \langle N_W \rangle \kappa_4(\Delta n) + 4 \langle \Delta n \rangle \kappa_3(\Delta n) \kappa_2(N_W)$ $+ 3 \kappa_2^2(\Delta n) \kappa_2(N_W) + 6 \langle \Delta n \rangle^2 \kappa_2(\Delta n) \kappa_3(N_W) + \langle \Delta n \rangle^4 \kappa_4(N_W).$

Additional terms appear from the event by event participant fluctuation

Δn : net-proton per N_W
 ΔN : net-proton

P. Braun-Munzinger, A. Rustamov, J. Stachel: NPA.2017.01.011



手法比較

Centrality bin width correction

- ・ モデルに依存しない。
- ・ Multiplicity 1 ビンの分解能で補正が頭打ちになる。
- ・ 補正が衝突エネルギーに依存する。

Volume fluctuation correction

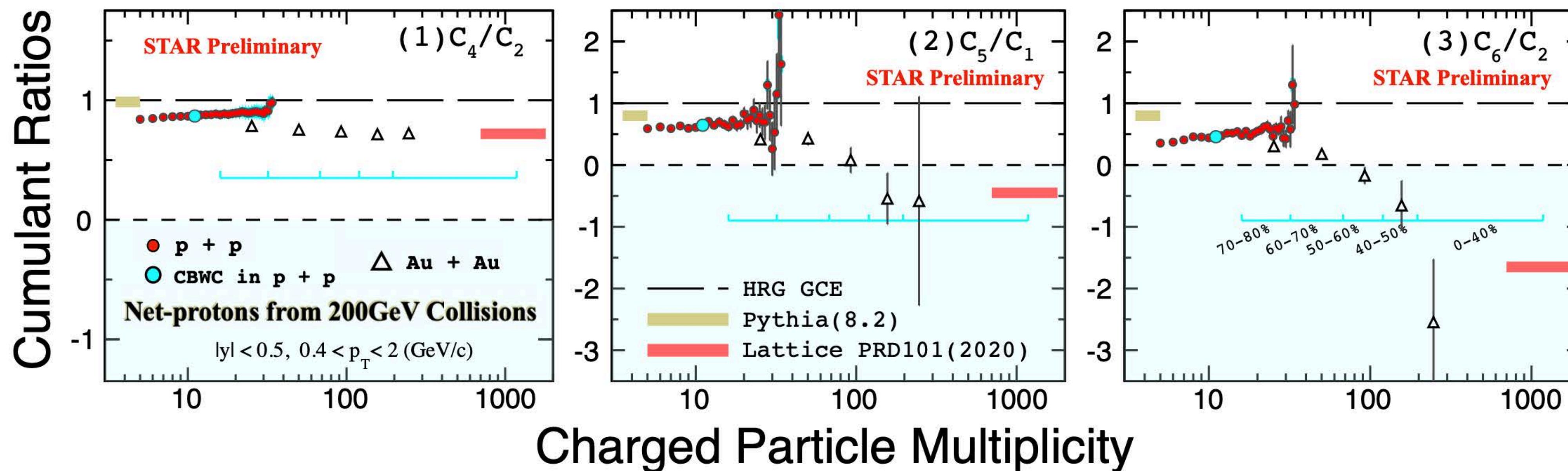
- ・ 独立な粒子生成モデル。
- ・ 初期ゆらぎ+粒子生成ゆらぎの 2 段階で仮定が必要。
- ・ これらの仮定のもと、体積ゆらぎを完全に除去できる。

問題点

- ・ たとえインパクトパラメータを直接測定できたとしても、解決する問題では無い。→ そもそも「初期体積」とは？
- ・ 現状では、 $(\text{体積ゆらぎ}) = (\text{初期ゆらぎ}) + (\text{粒子生成ゆらぎ})$ であり、これらを切り分けられないのが問題をさらに複雑にしている。
- ・ p+pとの比較で何かできないか？

$p+p$ vs $A+A$?

✓ $p+p$ と $Au+Au$ の差が体積ゆらぎだと仮定して、何か調べられないか？



- Only statistical errors are shown for $Au+Au$ results
- Efficiency is not corrected for x-axis

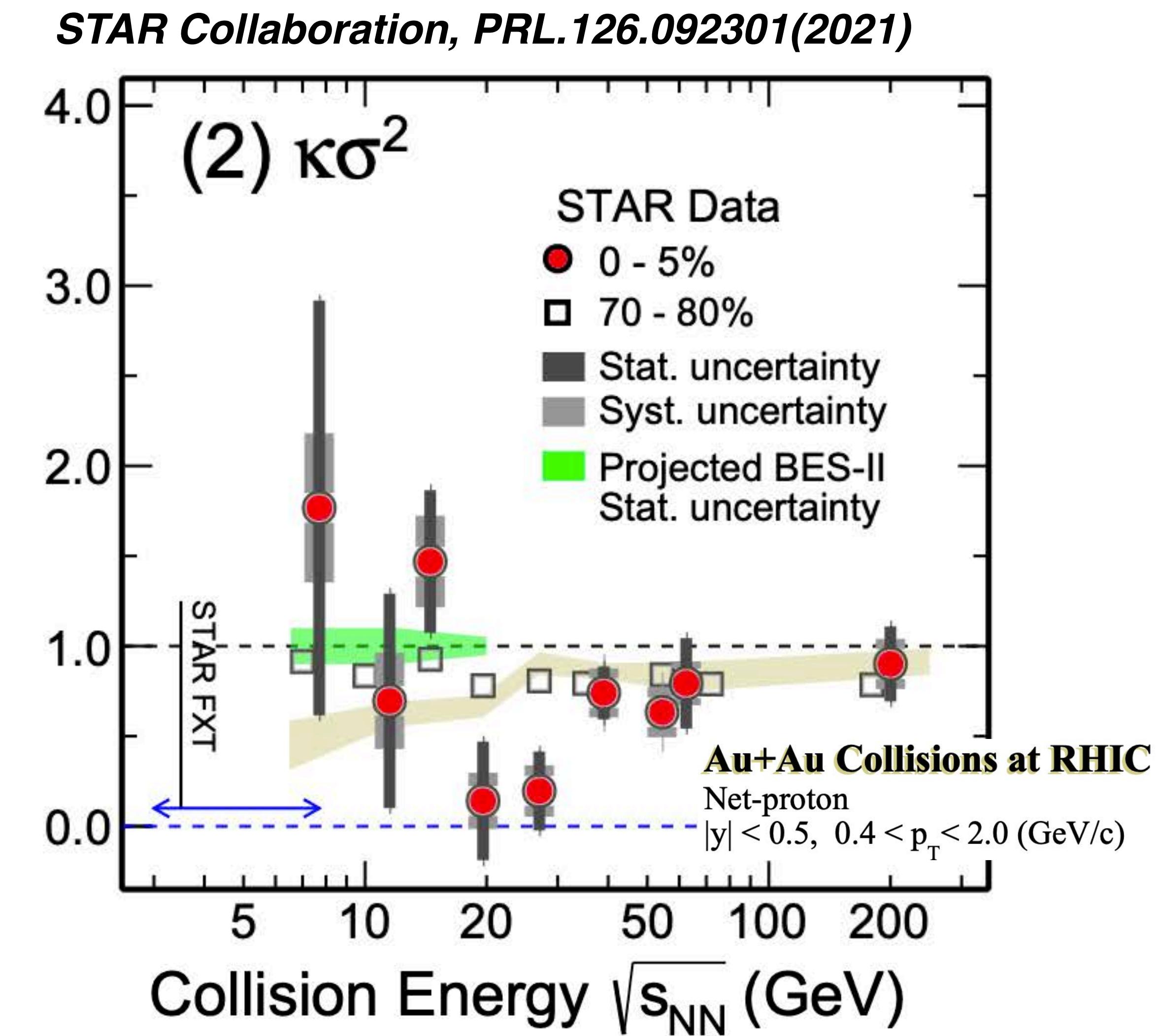
STAR Collaboration,
PRC.104.024902(2021)

LQCD : Phys. Rev. D 101, 074502 (2020)

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121882 (2021)

まとめ

- RHIC-STARにおける高次ゆらぎ測定で臨界点やクロスオーバーの兆候
- 今後（おそらく）3-4年でBES-II / FXTの結論が出る。
- 非単調な振る舞いが再確認された場合、その解釈は？
- 実験と比較ができる（体積ゆらぎを含む）動的モデルが必要。



Thank you for your attention